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# **Doppler seismology**

# **with a Michelson**

# **Interferometer**

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# Detection of p-modes

## Two methods:

- by photometry

The stellar oscillations generate temperature pulsations.

Repeated recording of the visible photometry of a star:

~ every min, over several days

Precision required:

$10^{-6}$

⇒ Measurements with time of photometric micro-variations

**Only feasible from space**

Space missions dedicated to asteroseismology:

MONS, MOST, COROT

# Detection of p-modes

## Two methods:

- **by spectroscopy**

The stellar oscillations generate surface motions.

Repeated recording of a high-resolution, visible spectrum of a star: ~ every min, over several days

⇒ Measurements with time of small Doppler shifts (after subtraction of the mean radial velocity)

## **Feasible from ground**

Need to set up a worldwide network of dedicated, échelle spectrometers



# Specifications of a ground-based

## grating spectrometer dedicated to asteroseismology:

- \* high spectral resolution  $\sim 10^5$
- \* broad spectral coverage *visible range*

## a fiber-fed cross-dispersed échelle spectrometer

- \* high stability *temperature control  $\sim 10^{-2} K$*
- \* spectral reference  $10^{-9}$
- \* minimum number of instruments for a network 3

# Specifications of a ground-based

grating spectrometer dedicated to asteroseismology:

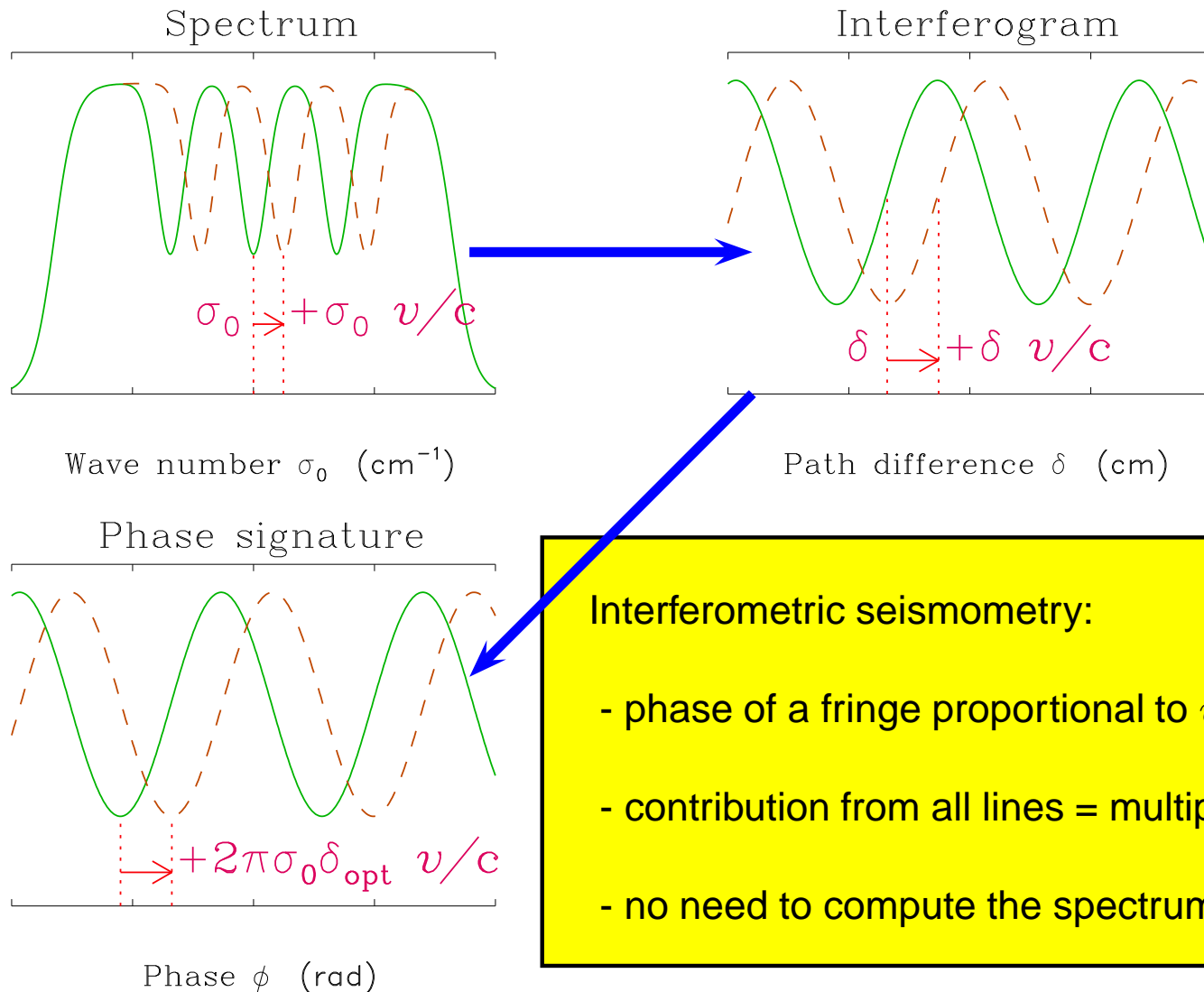
- \* high spectral resolution  $\sim 10^5$
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a fiber-fed cross-dispersed échelle spectrometer

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Another approach: with a Michelson interferometer

# Seismometry in the Fourier space



## Interferometric seismometry:

- phase of a fringe proportional to  $v$
- contribution from all lines = multiplex advantage
- no need to compute the spectrum

# What OPD to use?

Interferogram  $I(\delta)$  of a single absorption line at  $\sigma_0$ , depth  $A_0$ , visibility function  $\Gamma(\delta)$ :

$$I(\delta) = A_0 \Gamma(\delta) \cos 2\pi\sigma_0\delta$$

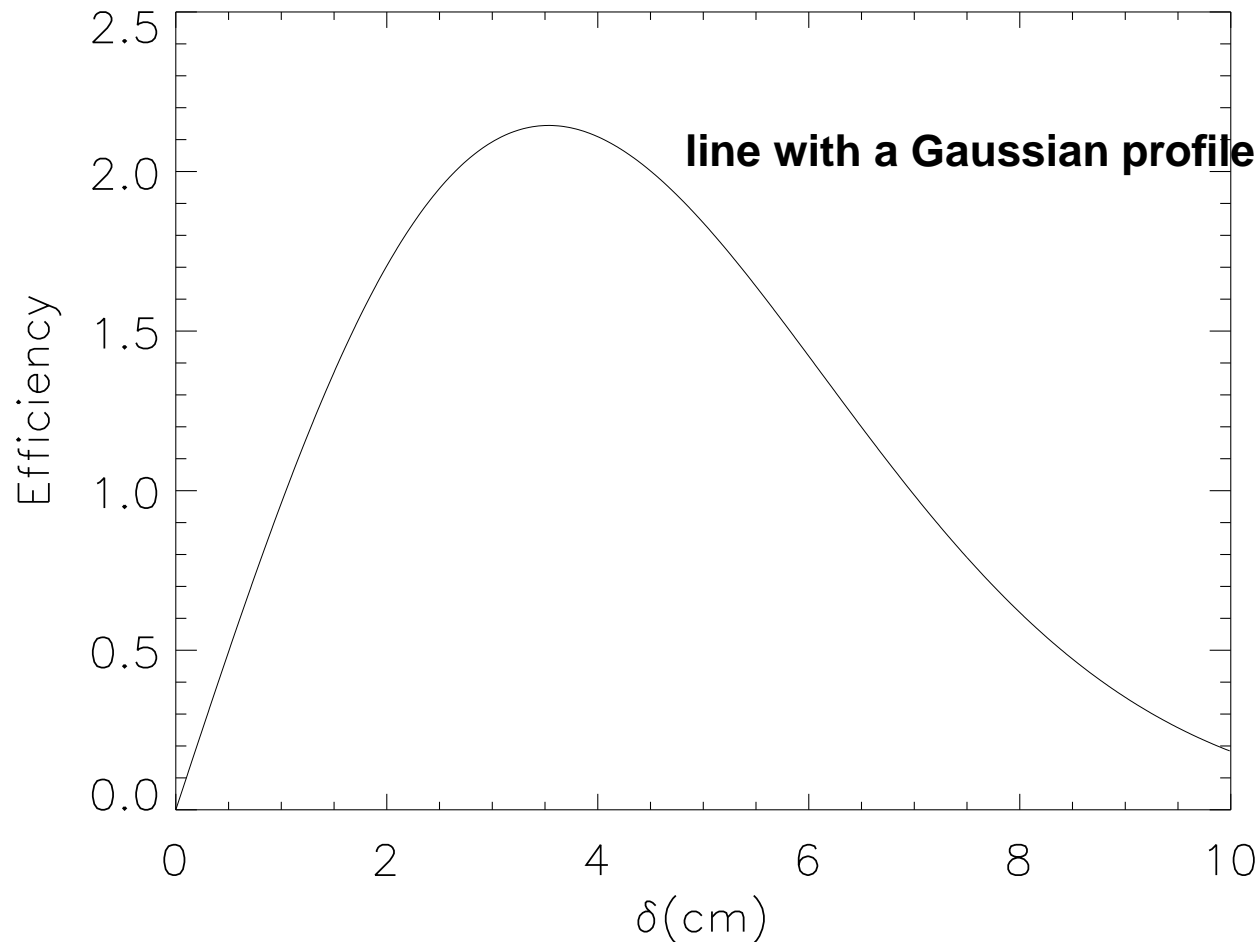
Modulation of the output signal by Doppler velocity  $dv$ :

$$dI(\delta) = 2\pi \frac{\sigma_0}{c} A_0 \delta \Gamma(\delta) \sin 2\pi\sigma_0\delta dv$$

Maximum signal at  $\delta_0 =$  zero crossing of a fringe (sinus = 1)

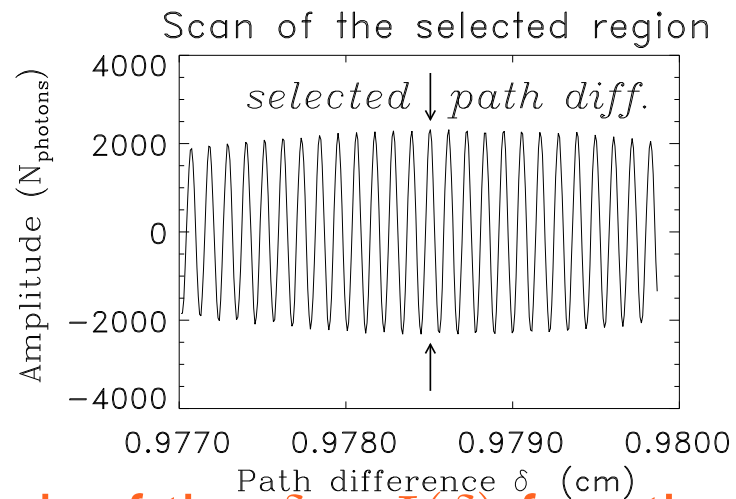
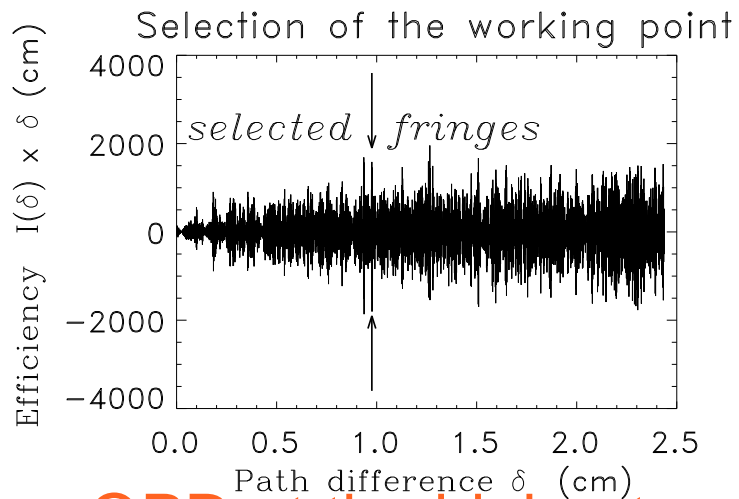
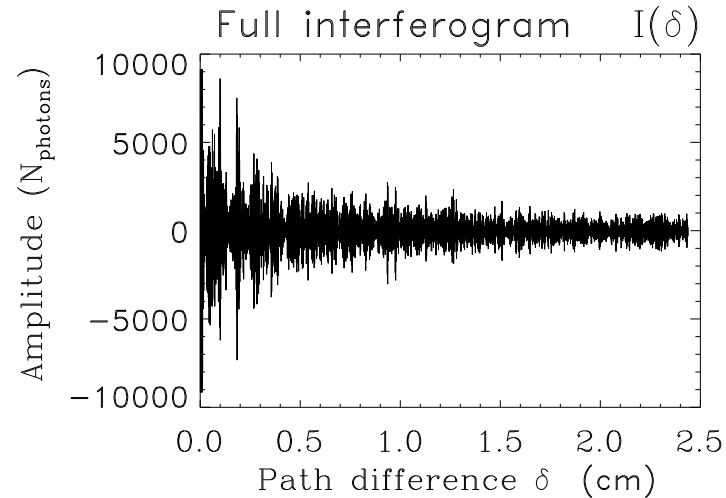
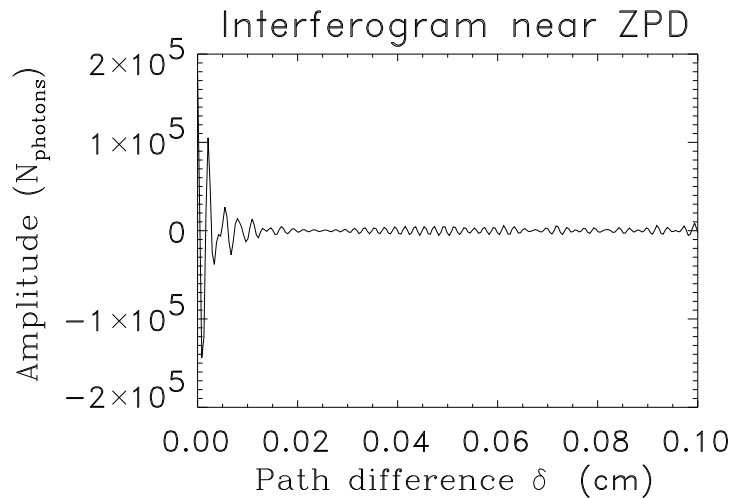
$$dI(\delta_0) = 2\pi \frac{\sigma_0}{c} A_0 \delta_0 \Gamma(\delta_0) dv$$

# What OPD to use?



$$\delta \times \Gamma(\delta) = \text{Efficiency fonction}$$

# Selection of the optimum OPD



**Optimum OPD at the highest peak of the  $\delta \times I(\delta)$  function**

From observations with the CFHT-FTS: (Mosser, Maillard et al. 2000, *Icarus*, 144, 104).

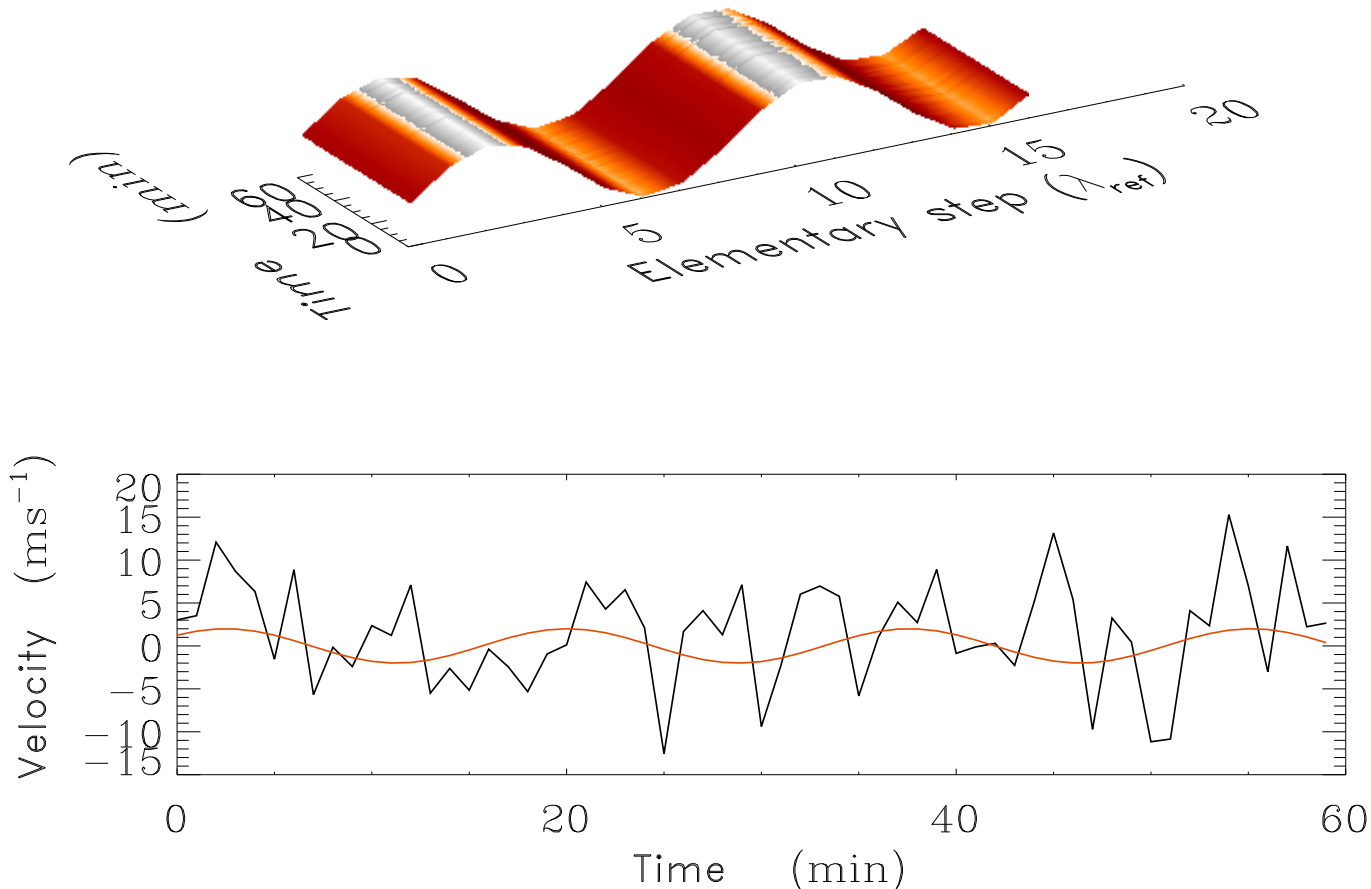
# Data acquisition (1)

**In principle**, simply with the interferometer blocked at  $\delta_{opt}$ , on a fringe zero-crossing, recording of the oscillation signal on the linear part of the selected fringe:

- Sensitive to the scintillation noise
- Sensitive to the transparency fluctuations
- Sensitive to the drifts of OPD (no velocity detection at the peak of a fringe)

**Solution: the recording of a phase signal**

# Data acquisition (2)



Continuous step-by-step scanning of the same fringe. Fitting of a sine wave through the recorded samples (minimum 3)



# Fundamental performances

## Simulations:

from synthetic spectra of dwarf stars:  $T_{eff}$  5000 to 7000K  
range: 380 - 680 nm (14700 - 26300  $cm^{-1}$  )

## Parameters:

V magnitude: 0 to 6  
rotational velocity  $v \sin i$ : 0 to 80  $km s^{-1}$

## Conditions:

2-m class telescope

photon noise-limited by the star flux

# MAIN RESULTS OF SIMULATIONS

- Comparison of two instrument configurations

- unique narrow-band spectral domain: two detectors

**FS1**

- multi-bandpass, by post-dispersion: array detectors

**FS $\lambda$**

*Photon noise-limited Doppler measurements with a Fourier Transform Seismometer I. Fundamental Performances*

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- Influence of the stellar *vsini*

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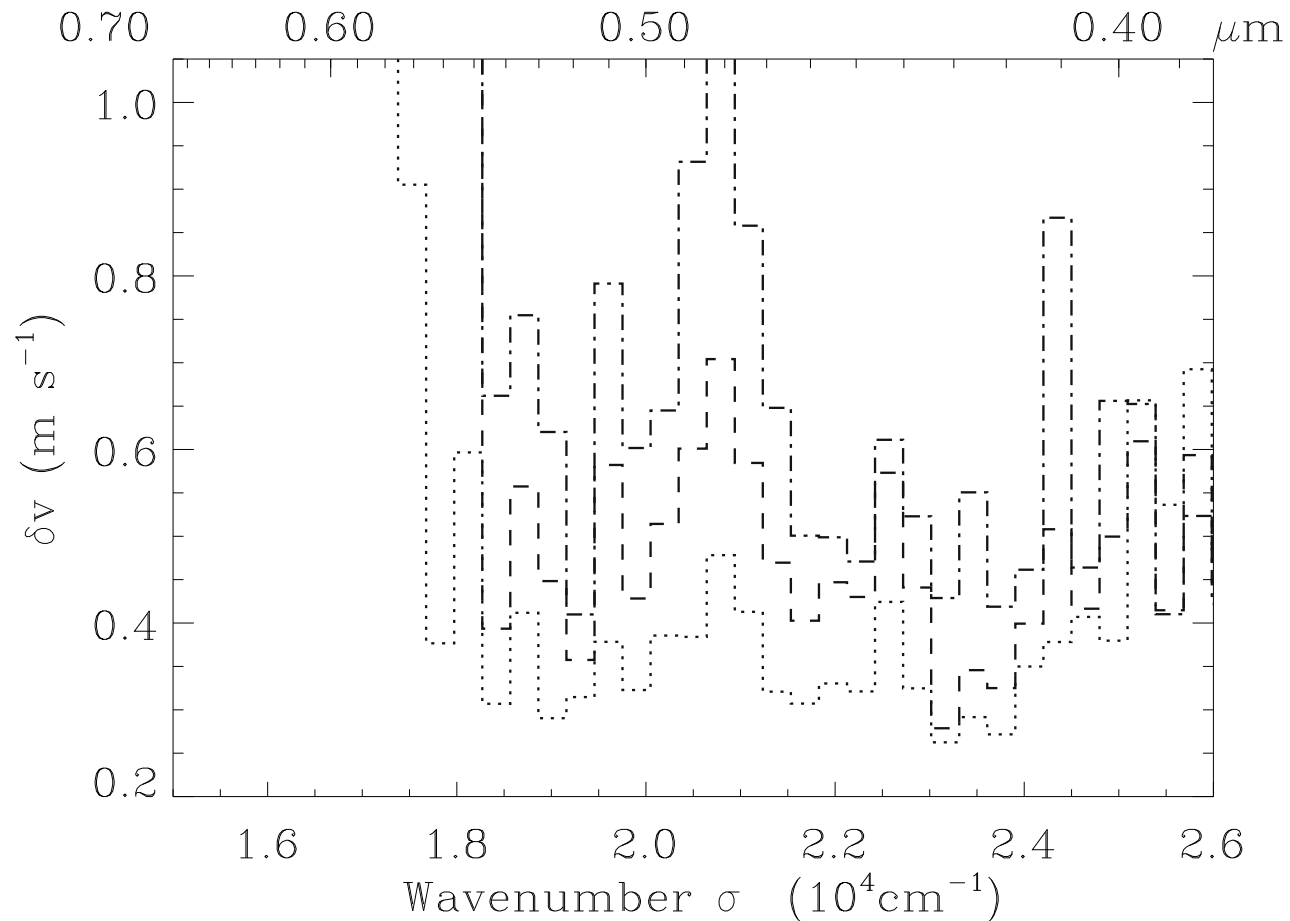
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- Comparison with a grating spectrometer

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# Spectral domain



Limit of velocity detection ( $\delta v$ ) due to photon noise as a function of the mean  $\lambda$  of the bandpass for three different dwarf stars (*dot* K2, *dash* F9, *dot-dash* F2).

Conclusion: **combined with CCD efficiency  $\Rightarrow$  400 – 560 nm**

# Comparison of the two configurations

– two detectors, unique narrow-band spectral domain

**FS1**

*Mode tested with the CFHT-FTS on stars and on Jupiter*

– array detectors, multi-bandpass by post-dispersion

**FS $\lambda$**

*Obtained by a low-resolution grating on each parallel, output beam of the interferometer (dual output)*

Target	mag	vsini km s <sup>-1</sup>	FS1 $\delta v$ cm s <sup>-1</sup>	FS $\lambda$ $\langle \delta v \rangle$ cm s <sup>-1</sup>	$\delta v_{\text{FS1}}/\delta v_{\text{FS}\lambda}$
F3V	6.3	12.3	51	17.1	<b>3.0</b>
F7V	5.0	11.2	23	7.6	<b>3.0</b>
G0V	4.8	10.9	19	5.7	<b>3.3</b>
G5V	5.4	10.8	18	5.7	<b>3.2</b>
K3V	6.2	12.5	17	5.7	<b>3.0</b>



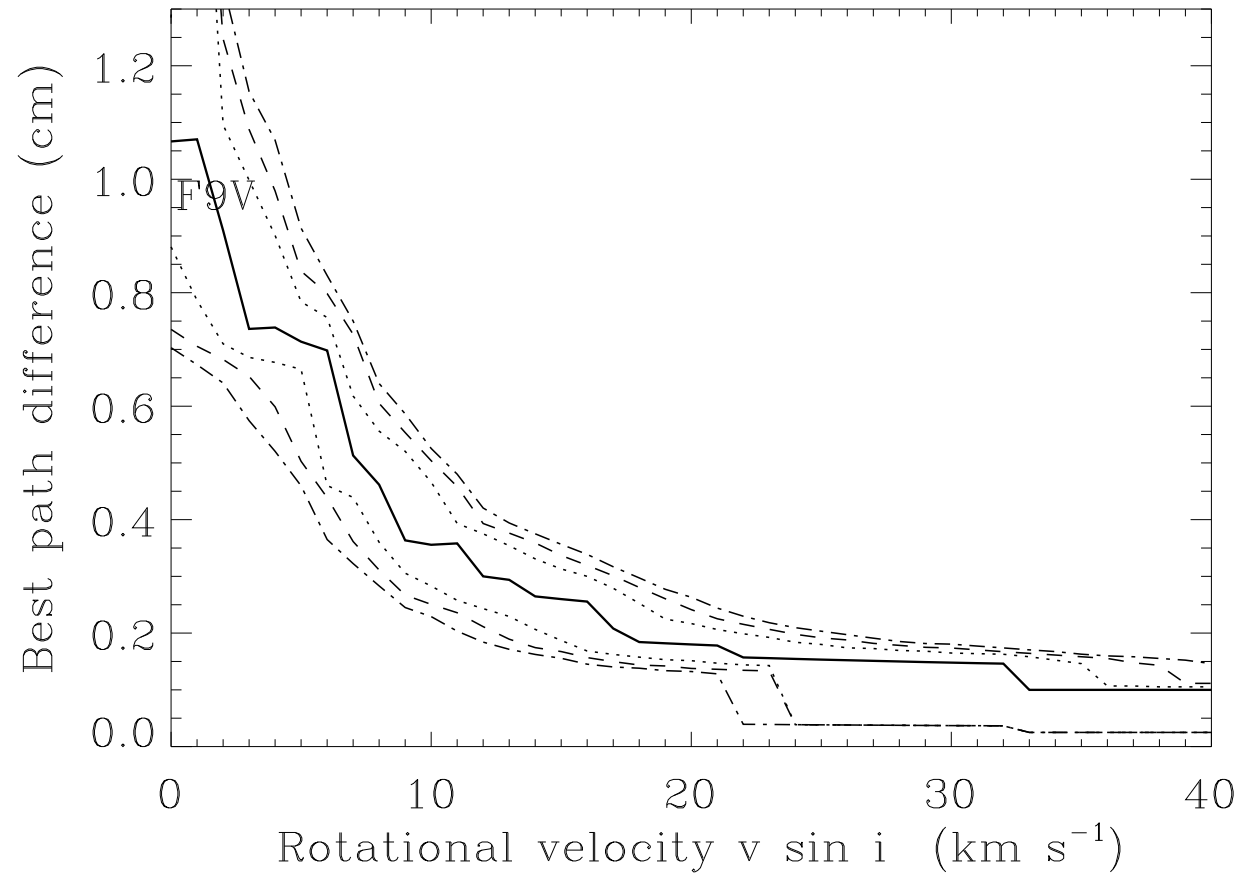
# Comparison of the two configurations

## Conclusions:

- Multi-bandpass better by a factor  $\sim 3$  than the mono-bandpass configuration.
- Gain by the coverage of a large spectral range, without increasing the photon noise.

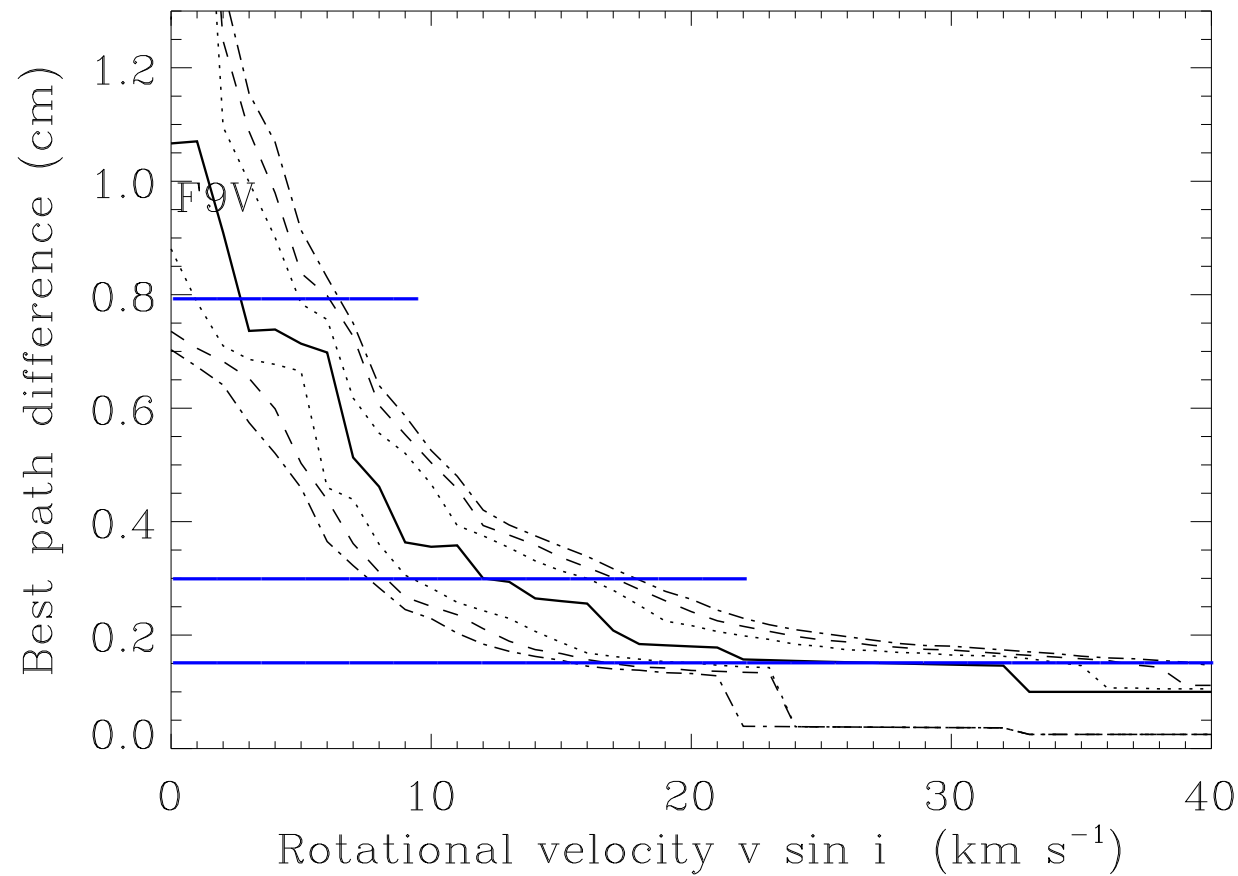
Necessary to adopt a montage with a post-disperser

# Optimum OPD versus $v \sin i$



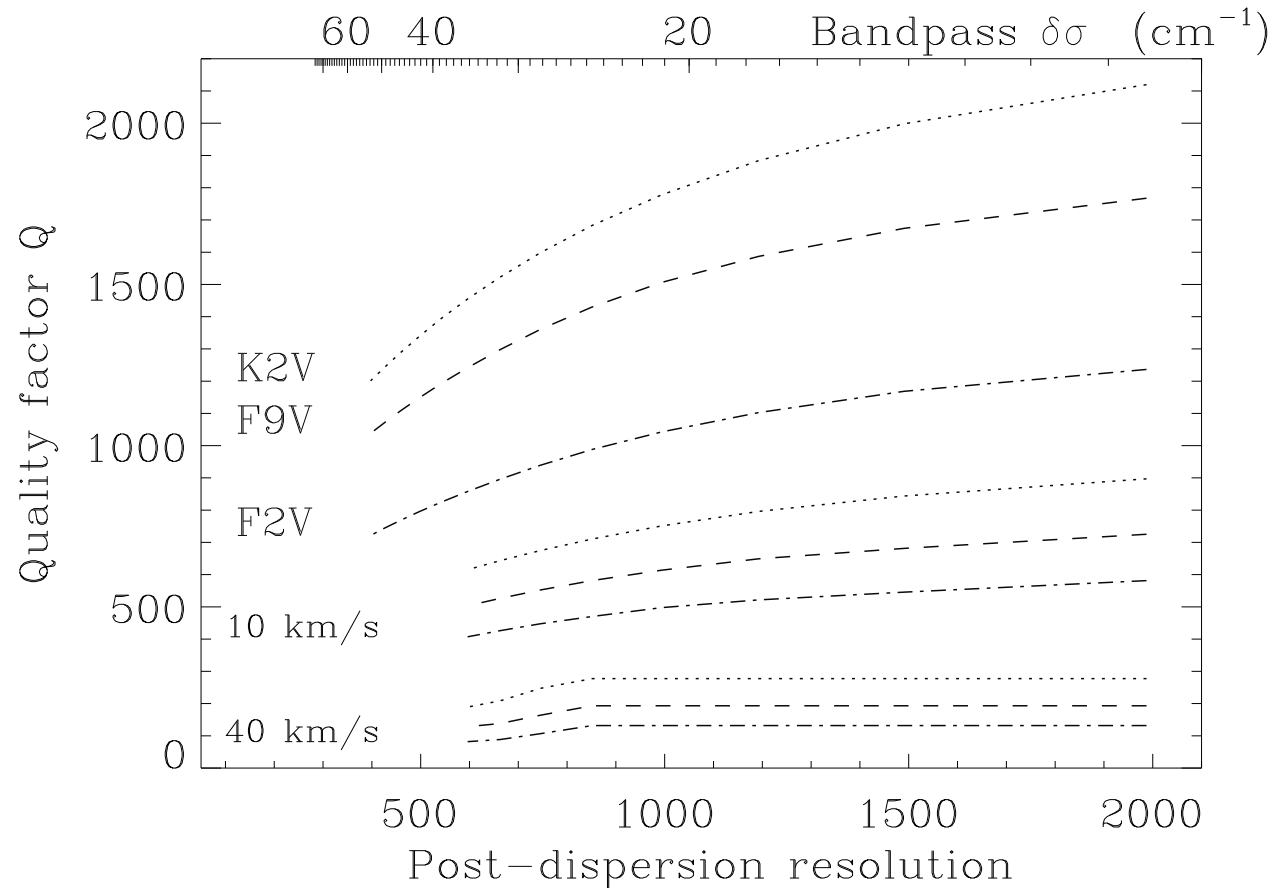
Best OPD as a function of  $v \sin i$  for a F9V star (—). Reduction of performances by 5% (dot), 10% (dash) and 15% (dot-dash).

# Optimum OPD versus $v \sin i$



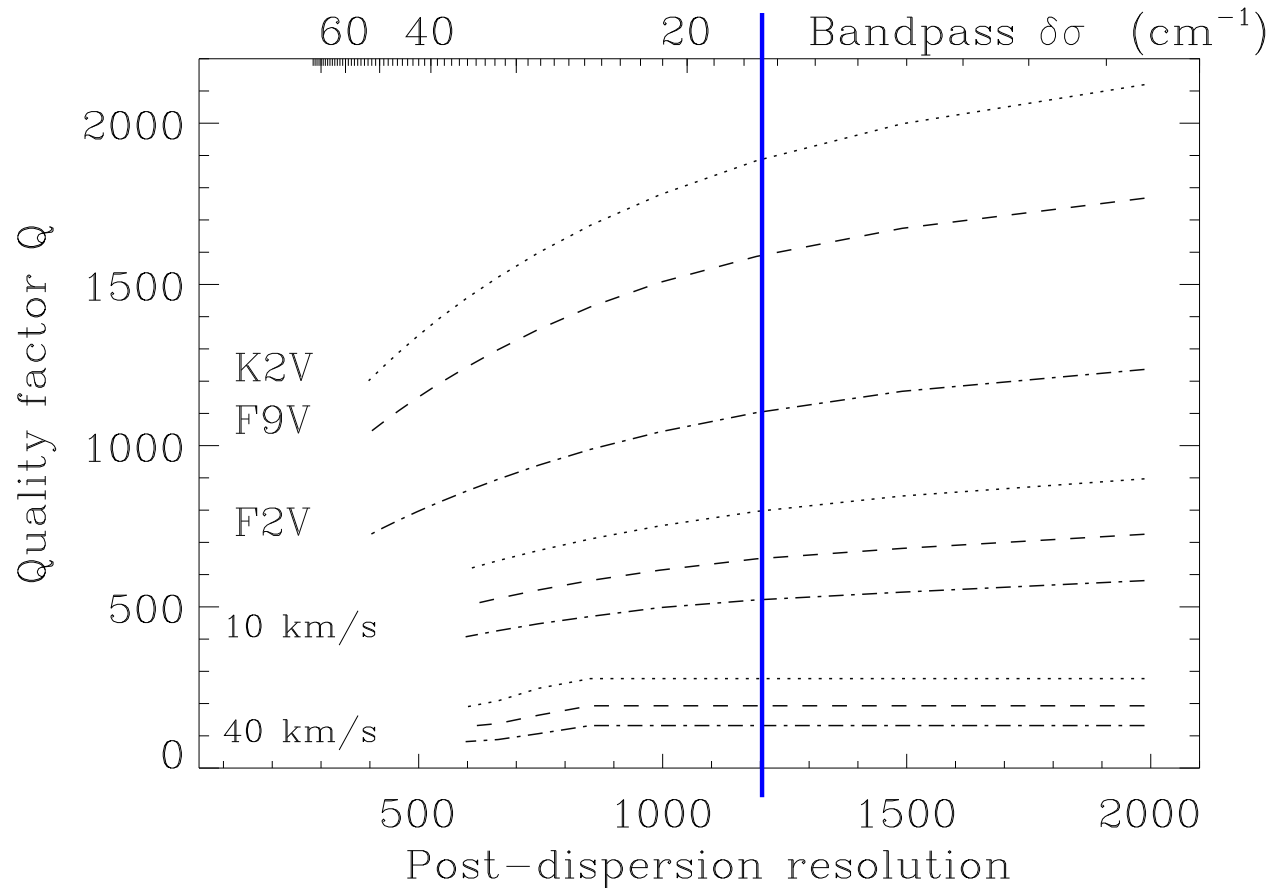
Conclusion: to cover the full range of  $v \sin i$   
3 OPDs needed, 0.8cm, 0.3cm, 0.15cm

# Influence of post-disperser resolution



Quality factor of the stellar spectrum as a function of the resolution of the post-disperser for  $v \sin i = 0, 10$  and  $40 \text{ km s}^{-1}$  and 3 stellar types K2V (*dot*), F9V (*dash*) and F2V (*dash-dot*)

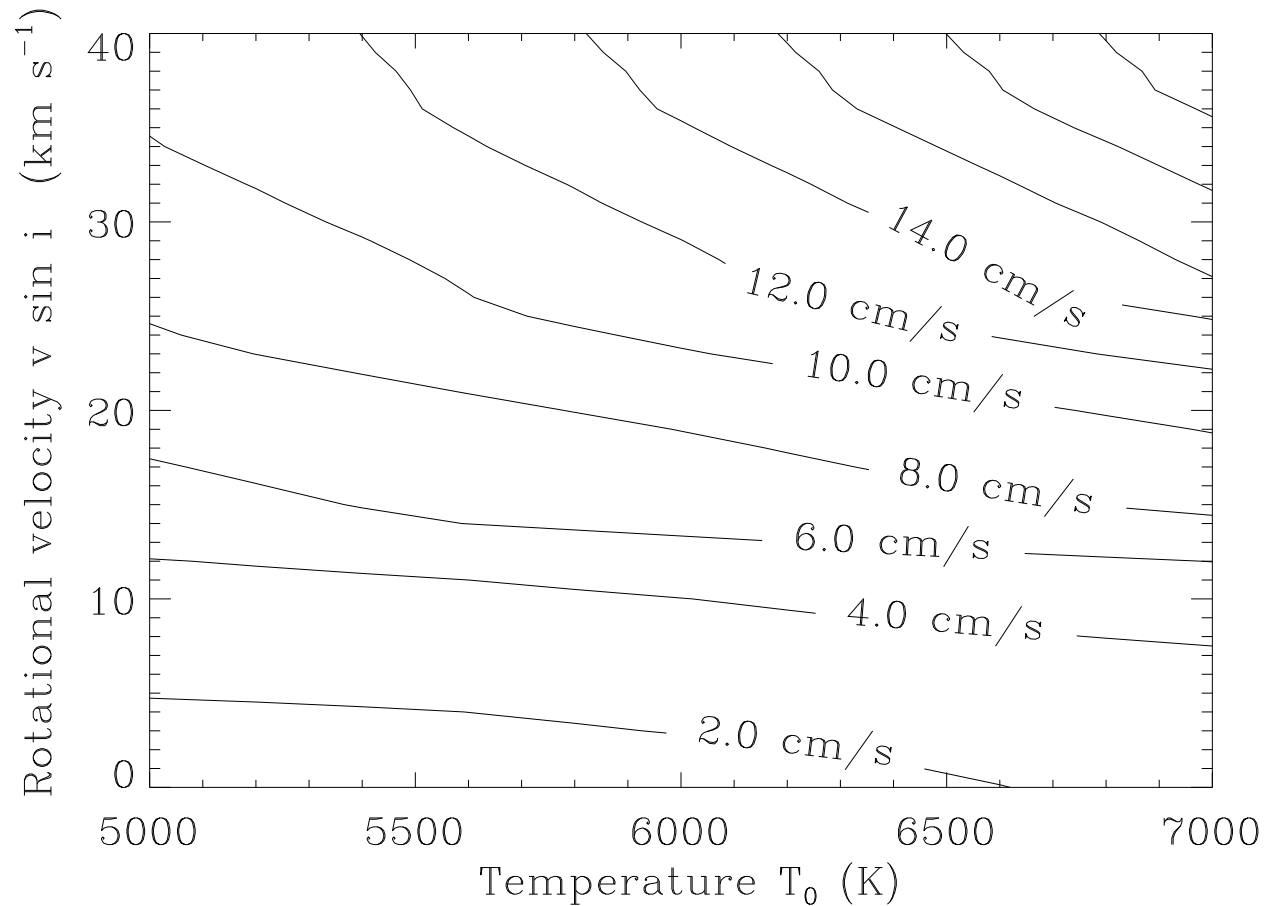
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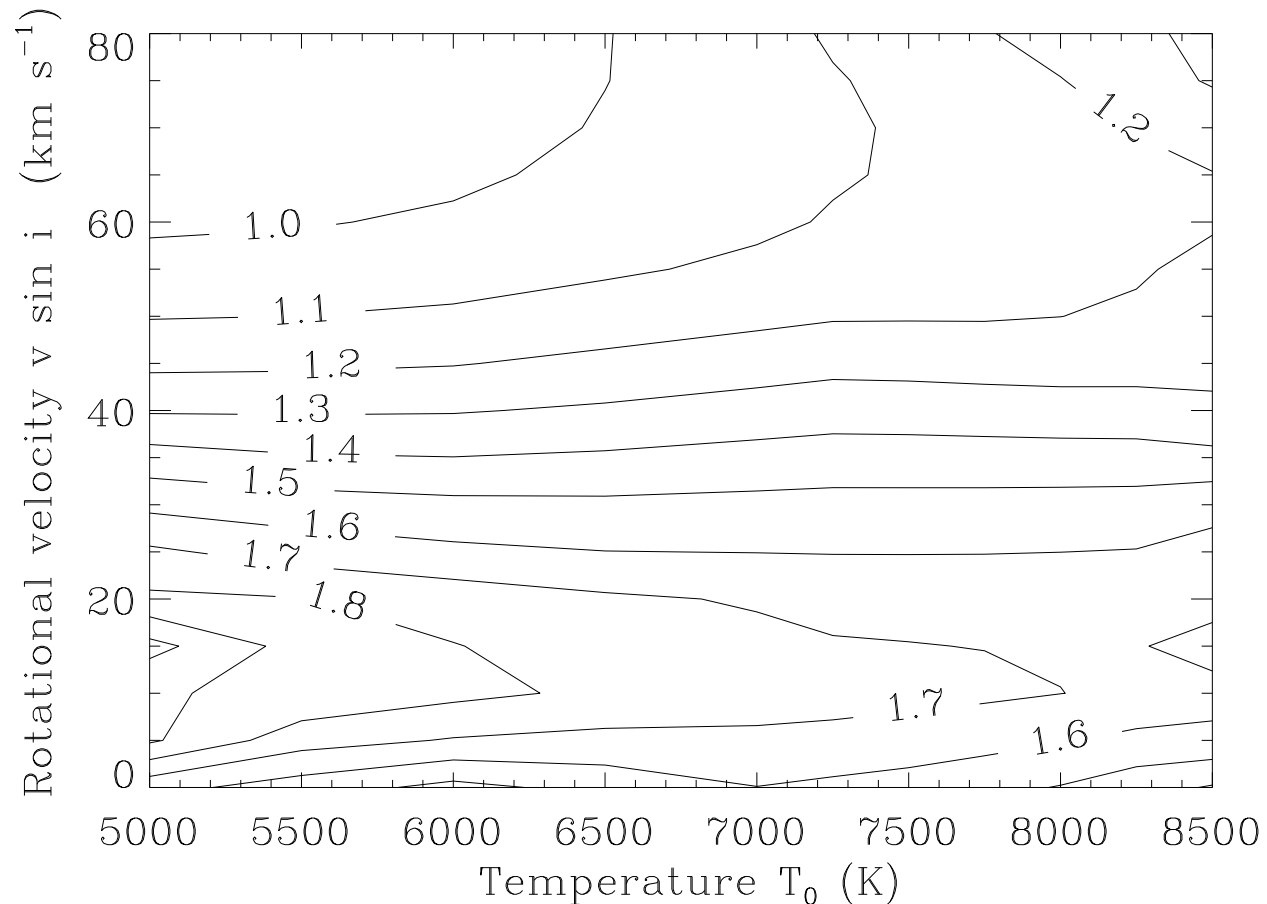
Conclusion: resolution of post-disperser  $\Rightarrow$  1200

# Influence of $v \sin i$ on performances



$\delta v_{rms}$  as a function of effective stellar  $T_0$  and  $v \sin i$

# Comparison with a grating spect. (GS)



Ratio of the performances of  $\delta v_{\text{FTS}}/\delta v_{\text{GS}}$  as a function of  $T_0$  and  $v \sin i$   
on the same star, for the same observing conditions ( $R_{\text{GS}} = 88\,000$ ,  $R_{\text{post-disp}} = 1200$ ).

**Conclusion: FTS < GS by a factor between 1 and 1.8**

# The SIAMOIS concept

stands in French for:

**S**ismomètre **I**nterferentiel **A** **M**esurer *les*  
**O**scillations **S**tellaires



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All the results of the simulations used to define the characteristics of a FTS dedicated to asteroseismology

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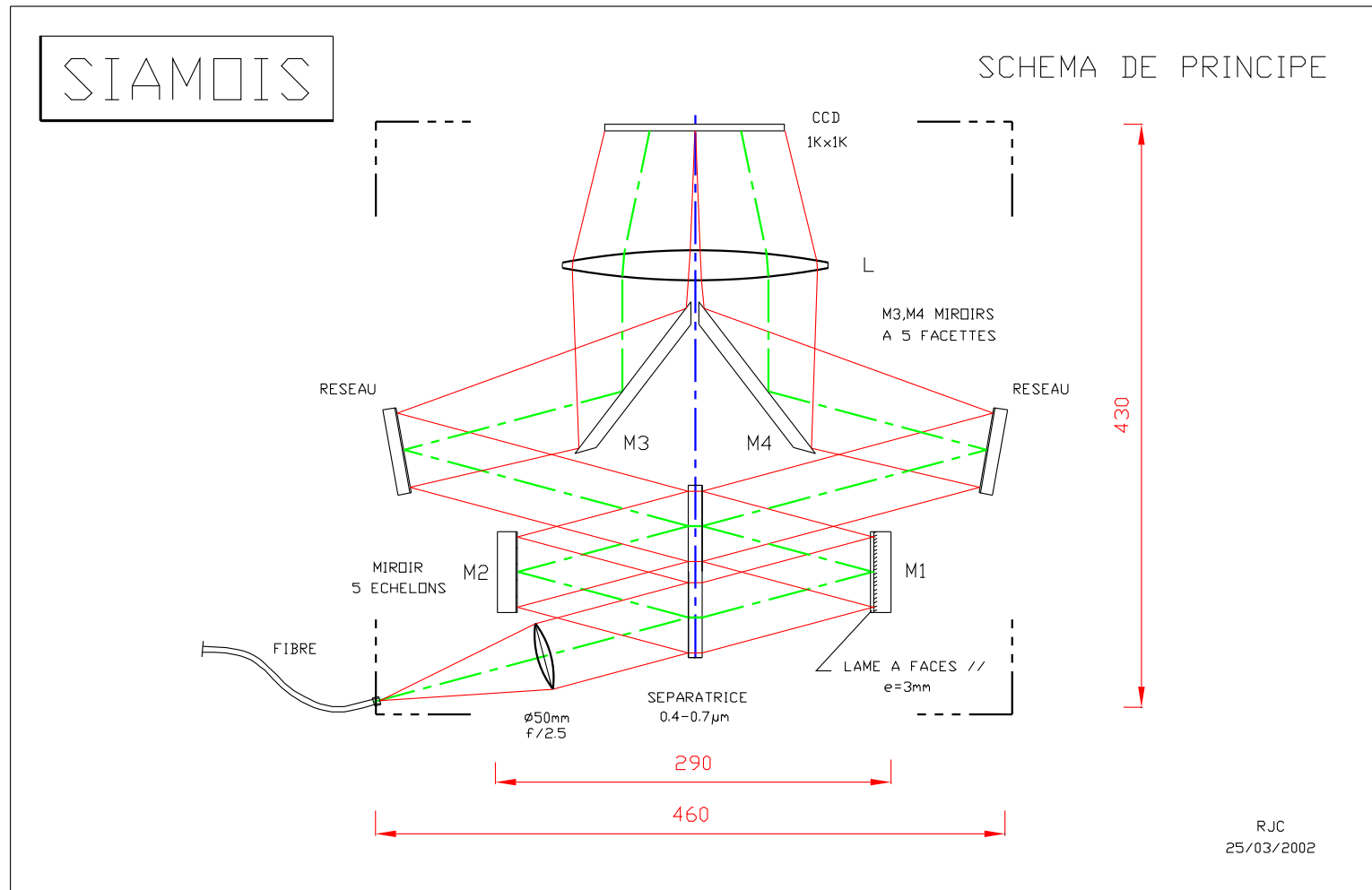
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- high stability *no adjustment*
- simplicity *no servo-system, no electronics control*
- no scanning *all samples recorded simultaneously*
- modest size *small path difference*
- adaptable to Dome C conditions

**A monolithic interferometer, with no moving parts**

# Schematic diagram of SIAMOIS

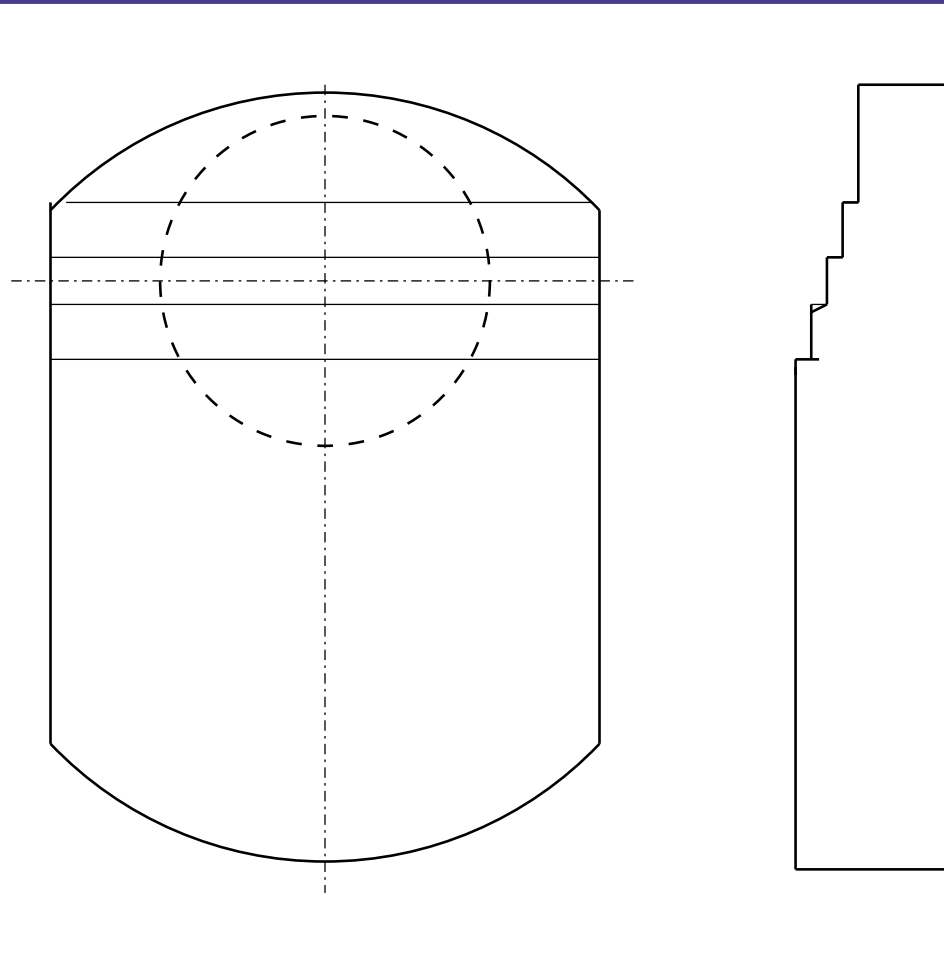


The concept is based on a compact Mach-Zehnder interferometer:  
a single beamsplitter and two output beams

# Elements of the SIAMOIS design

- feeding of the instrument from the telescope by optical fiber  
stabilization of star image
- fixed OPD  
a parallel plate in one arm
- sampling of a fringe  
in 5 points
- one mirror of the interferometer  
5 steps of depth =  $\lambda_{mean}/10$

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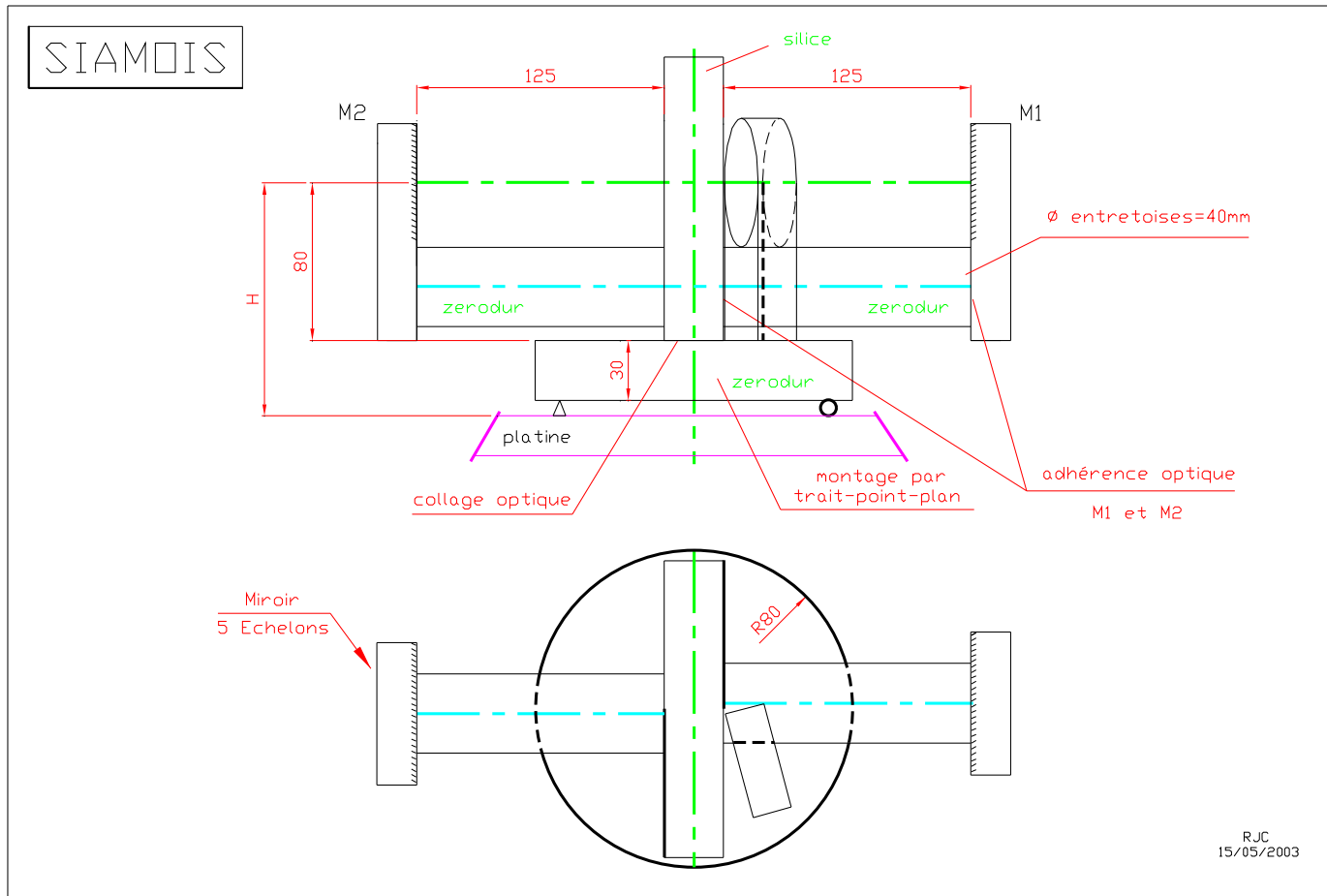


Profile of the 5-step mirror. Depth of each step  $\lambda_{mean}/10$ . Width of each portion of the beam calculated to obtain about 5 equal intensity channels

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in 5 points
- one mirror of the interferometer  
5 steps of depth =  $\lambda_{mean}/10$
- two beam slicer mirrors  
5 channels
- post-dispersion by two low-resolution gratings  
R=1200
- a blue CCD  
1024×256

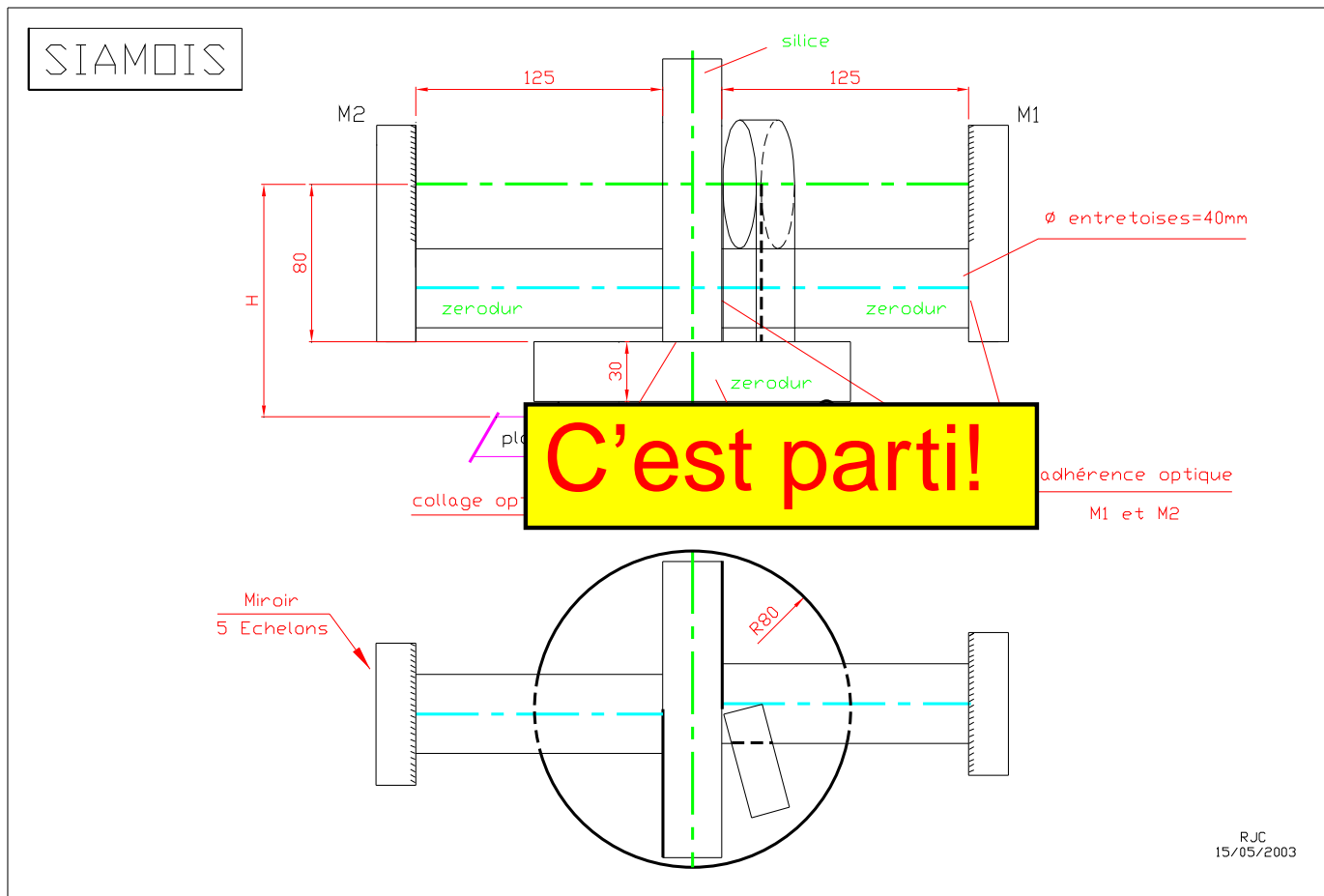
# A monolithic interferometer



**Fig. 3** Ensemble interférométrique monolithique avec miroirs et séparatrice.  
Montage de la lame à faces parallèles.

The two mirrors and the beamsplitter make a single optical block with two Zerodur cylindrical spacers optically contacted. The block is bonded on the optical bench.

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